

Mathematics

On One Singular Integral Equation Arising from the Radiative Transfer Theory

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ABSTRACT. In the class of Hölder functions we give the necessary and sufficient condition for solvability of the two-dimensional integral equation having singularity with respect to one variable. Such equations often arise from the radiative transfer theory. Finding a solution is reduced to solving a one dimensional Fredholm integral equation of the second kind. © 2015 Bull. Georg. Natl. Acad. Sci.

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We present a method for solving singular integral equation occurring while investigating many important problems of neutron transport theory [1,2]. This equation has the form

$$(Lu)(x,t) \equiv \int_{a-1}^{b+1} \frac{sk(x,y)}{s-t} u(y,s) ds dy + u(x,t) + \int_{a-1}^{b+1} \frac{tk(x,y)}{s-t} u(y,t) ds dy = f(x,t), \quad (1)$$

$$x \in [a,b], t \in (-1,+1)$$

where $k(x,y)$ is the real valued continuous function, the right part $f(x,t)$ is the real valued function satisfying H^* condition [3] with respect to t . Further, we shall denote the class of such functions by D^* . We look for solutions $u(x,t) \in D^*$. To our opinion it appears to be more natural and appropriate in certain applications of equations. Note that in this paper we generalize the result from [6], therefore formalism sometimes used here is similar to the formalism, which is used in [6].

Let us introduce the linear integral operator defined by formula:

$$(\Omega_z g_z)(x,t) \equiv g_z(x,t) + z \int_{a-1}^{b+1} \frac{k(x,y)}{s-z} g_z(y,s) ds dy,$$

$$z \notin [-1,+1], x \in [a,b], t \in (-1,+1),$$

where the parameter z is any point on the plane. The boundary properties of the introduced operator are of great importance for further investigation of the operator L . The operator Ω_z operating on any continuous function $g_z(x, t)$ piecewise holomorphic in z with a cut on the real axis $[-1, +1]$ and satisfying the H^* condition in t will define a piecewise holomorphic function with a cut on $[-1, +1]$. Using the Plemelj-Sokhotski formulas [3], we can calculate the limiting values of $g_z(x, t)$ as

$$\Omega_{\zeta t}^{\pm} g_{\zeta}^{\pm}(x, t) = g_{\zeta}^{\pm}(x, t) + \zeta \int_{a-1}^{b+1} \frac{k(x, y)}{s - \zeta} g_{\zeta}^{\pm}(y, t) ds dy \pm i\pi \int_a^b k(x, y) g_{\zeta}^{\pm}(y, t) dy, \tag{2}$$

$$\zeta \in (-1, +1), \quad x \in [a, b], \quad t \in (-1, +1).$$

Denote by \mathfrak{R} the set of values of z , for which the homogeneous equation

$$\Omega_z u = 0 \tag{3}$$

admits non-zero solution. Such values of z are called eigenvalues of Ω_z . In as much as the kernel of the operator Ω is a piecewise analytic function in z , satisfying the H condition [3] and bounded as $z \rightarrow \infty$, from Tamarkin's Theorem [4] it follows that the set \mathfrak{R} is no more than countable in the plane with a cut on $[-1, +1]$.

It is not difficult to prove that:

I_1 . If $g_{z_k}(x, t)$ is a solution of Eq.(3), then the function

$$\tau_{z_k}(x, t) = \frac{z_k g_{z_k}(x, t)}{z_k - t},$$

where $z_k \in \mathfrak{R}$ is a solution of the equation

$$(z - t)\tau_z(x, t) - z \int_{a-1}^{b+1} k(x, y)\tau_z(y, s) ds dy = 0,$$

$$x \in [a, b], \quad t \in (-1, +1),$$

when $z = z_k$ and vice versa. In addition $\mathfrak{R} \cap [-1, +1] = \emptyset$.

I_2 . If $z = z_k$ is an eigenvalue of multiplicity r_0 for the operator Ω_z , then z_k also is an eigenvalue of multiplicity r_0 for the operator

$$(\Omega_z^* q_z)(x, t) \equiv q_z(x, t) + z \int_{a-1}^{b+1} \frac{k(y, x)}{s - z} q_z(y, s) ds dy$$

$$z \notin [-1, +1], \quad x \in [a, b], \quad t \in (-1, +1),$$

and vice versa.

I_3 . If $z \neq z'$, then

$$\int_{a-1}^{b+1} t \tau_z^*(x, t) \tau_{z'}(x, t) dt dx = 0,$$

where $\tau_z^*(x, t)$ is a solution of the equation

$$(z-t)\tau_z^*(x, t) - z \int_{a-1}^{b+1} k(y, x)\tau_z^*(y, s)dsdy = 0,$$

$$x \in [a, b], t \in (-1, +1).$$

In the sequel we consider the case when \mathfrak{R} is the finite set. The functions $\tau_{z_k}^*(x, t)$ and $\tau_{z_k}(x, t)$ $z_k \in \mathfrak{R}$, are called eigenfunctions of the kernel $k(y, x)$ and $k(x, y)$, respectively. The basic result for L is the following theorem.

Let $f(x, t) \in D^*$.

Theorem. For the singular integral equation

$$Lu = f \tag{4}$$

to have a solution in the class D^* , it is necessary and sufficient that the function f satisfy the conditions

$$\int_{a-1}^{b+1} t \tau_{z_k}^*(x, t) f(x, t) dt dx = 0, \quad z_k \in \mathfrak{R}. \tag{5}$$

Proof. Proof of the necessity. Let us introduce into consideration the piecewise holomorphic function

$$\psi_z(x) = \frac{1}{2\pi i} \int_{a-1}^{b+1} \frac{sk(x, y)}{s-z} u(y, s) ds dy,$$

where z is an arbitrary point on the plane. This function possesses the following properties:

(P)₁. In the plane with the cut $[-1, +1]$, it is piecewise holomorphic with respect to z .

(P)₂. As $z \rightarrow \infty$ it vanishes uniformly in x .

(P)₃. By using the Plemelj-Sokhotski formulas

$$\psi_t^\pm(x) = \frac{1}{2\pi i} \int_{a-1}^{b+1} \frac{sk(x, y)}{s-t} u(y, s) ds dy \pm \frac{t}{2} \int_a^b k(x, y) u(y, t) dy$$

$$x \in [a, b], t \in (-1, +1).$$

Combining the above equalities with (2), in view of (1) we obtain that the function $\psi_z(x)$ will be solution of the following boundary value problem:

$$(\Omega_t^+ \psi_t^+)(x) - (\Omega_t^- \psi_t^-)(x) = t \int_a^b k(x, y) f(y, t) dy,$$

$$x \in [a, b], t \in (-1, +1).$$

Consequently, taking the Plemelj-Sokhotski formulas into account, we can write

$$(\Omega_z \psi_z)(x) = \frac{1}{2\pi i} \int_{a-1}^{b+1} \frac{sk(x, y)}{s-z} f(y, s) ds dy,$$

$$x \in [a, b], t \in (-1, +1).$$

For this integral equation to define the analytic function $\psi_z(x)$ in a plane with a cut $[-1, +1]$, it is necessary and sufficient that its right-hand side satisfy the conditions

$$\int_{a-1}^{b+1} \int_{a-1}^{b+1} \tau_{z_k}^*(x, t) \int_{a-1}^{b+1} \frac{sk(x, y)}{s - z_k} f(y, s) ds dy dt dx = 0.$$

After simplification we get equality (5). The necessity is proved.

Proof of the Sufficiency. Let the function $f(x, t)$ satisfy the conditions (5).

By virtue of Tamarkin's theorem [5] it follows that this solution possesses the properties:

(R_1). It is piecewise holomorphic with respect to z .

(R_2). As $z \rightarrow \infty$ it vanishes uniformly in x .

(R_3). It may be represented in the form:

$$\psi_z(x) = \frac{1}{2\pi i} \int_{-1}^{+1} \frac{sm_s(x)}{s - z} ds, \\ z \notin [-1, +1], \quad x \in [a, b], \quad t \in (-1, +1),$$

where $m_z(x)$ is uniquely determined function.

In view of the Plemelj-Sokhotski formulas, the following equality is valid in the limiting values

$$\psi_t^+(x) + \psi_t^-(x) = \frac{1}{\pi i} \int_{-1}^{+1} \frac{\psi_s^+(x) - \psi_s^-(x)}{s - t} ds \tag{6} \\ x \in [a, b], \quad t \in (-1, +1).$$

Using (2), from (6) we are able to write

$$\overline{\psi}_t(x) + t \int_{a-1}^{b+1} \frac{sk(x, y)}{s - t} \overline{\psi}_t(y) ds dy + \int_{a-1}^{b+1} \frac{sk(x, y)}{s - t} \overline{\psi}_s(y) ds dy = \int_a^b k(x, y) f(y, t) dy, \tag{7}$$

where $\overline{\psi}_t(x) = \psi_t^+(x) - \psi_t^-(x)$.

Consider now the following equation

$$u(x, y) + \int_{a-1}^{b+1} \int_{a-1}^{b+1} \frac{tk(x, y)}{s - t} ds u(y, t) dy = \int_{-1}^{+1} \frac{s\overline{\psi}_s(x)}{t - s} ds + f(x, t), \tag{8} \\ x \in [a, b], \quad t \in (-1, +1),$$

which has a unique solution.

Denote by

$$p(x, t) = \overline{\psi}_t(x) - \int_a^b k(x, y) u(y, t) dy, \\ x \in [a, b], \quad t \in (-1, +1).$$

In view of (7) from (8) it follows that

$$p(x, t) + \int_{a-1}^{b+1} \int_{a-1}^{b+1} \frac{tk(x, y)}{s - t} p(y, t) ds dy = 0.$$

But Ω_z has no eigenvalues on $[-1, +1]$; consequently, we have

$$\bar{\psi}_t(x) = \int_a^b k(x, y)u(y, t)dy$$

Taking into account the last assertion, from (8) we obtain

$$u(x, t) + \int_{a-1}^{b+1} \frac{tk(x, y)}{s-t} u(y, t)dsdy = \int_{a-1}^{b+1} \frac{sk(x, y)}{t-s} u(y, s)dsdy + f(x, t),$$

$$x \in [a, b], t \in (-1, +1),$$

which means that (4) holds and the proof is complete.

Now, the aim is to study the singular operator L more deeply. We wish to find the inverse operator of L . For this purpose, in the class D^* we need to introduce the following singular operator

$$(Su)(x, t) \equiv \int_{a-1}^{b+1} \frac{tk(x, y)}{t-s} u(y, s)dsdy + u(x, t) + \int_{a-1}^{b+1} \frac{tk(x, y)}{s-t} u(y, t)dsdy,$$

$$x \in [a, b], t \in (-1, +1).$$

We shall note the following property of the introduced operator. For any two functions u and v from D^*

$$\int_{a-1}^{b+1} \int_{a-1}^{b+1} uSvdt dx = \int_{a-1}^{b+1} \int_{a-1}^{b+1} vLudt dx.$$

Consequently, if equation (1) has a solution, then necessarily

$$\int_{a-1}^{b+1} \int_{a-1}^{b+1} vfdt dx = 0,$$

where v is any solution of the homogeneous equation

$$Sv = 0.$$

In our further consideration we will need also the following identity:

$$\begin{aligned} & \int_{a-1}^{b+1} \int_{a-1}^{b+1} \frac{sk(x, y)}{t-s} \int_{a-1}^{b+1} \int_{a-1}^{b+1} \frac{s'k(y, y')}{s'-s} u(y', s')ds'dy'dsdy \\ &= \int_{a-1}^{b+1} \int_{a-1}^{b+1} \frac{u(y', s')}{s'-t} \left(\int_{a-1}^{b+1} \int_{a-1}^{b+1} \frac{sk(x, y)s'k(y, y')}{t-s} dsdy - \int_{a-1}^{b+1} \int_{a-1}^{b+1} \frac{sk(x, y)s'k(y, y')}{s'-s} dsdy \right) ds'dy' \\ & \quad + \pi^2 t^2 \int_a^b \int_a^b k(x, y')k(y, y')u(y', t)dy'dy', \end{aligned}$$

$$x \in [a, b], t \in (-1, +1).$$

This identity can be obtained by using the Bertrand-Poincare formula [3].

We can show the important property of the operator K introduced below:

$$(Ku)(x, t) = u(x, t) + \int_a^b r(x, y, t)u(y, t)dy,$$

$$x \in [a, b], t \in (-1, +1),$$

where

$$r(x, y, t) = \pi^2 t^2 \int_a^b k(x, y')k(y', y)dy' - \int_{-1}^{+1} \frac{tk(x, y)}{t-s} ds - \int_{-1}^{+1} \frac{tk(y, x)}{t-s} ds + \int_{a-1}^{b+1} \frac{tk(x, y')}{t-s} ds \int_{-1}^{+1} \frac{tk(y' y)}{t-s'} ds' dy'.$$

It is easily seen that if $k(x, y) = k(y, x)$, then the following decomposition exists

$$Ku = (B + i\pi C)(B - i\pi C)u,$$

where

$$(Bu)(x, t) \equiv u(x, t) + \int_{a-1}^{b+1} \frac{tk(x, y)}{t-s} u(y, t) ds dy$$

$$(Cu)(x, t) \equiv t \int_a^b k(x, y)u(y, t) dt.$$

In general case also, the similar decomposition is true. It follows that one-dimensional regular integral equation

$$h(x, y, t) + \int_a^b r(y', y, t)h(x, y', t)dy' = r(x, y, t), \quad x, y \in [a, b], \quad t \in (-1, +1)$$

admits unique solution. Now we can define the following integral operator

$$t(Tu)(x, t) = (S\dot{u})(x, t) + \int_a^b h(x, y, t)(S\dot{u})(y, t)dy,$$

$$x \in [a, b], \quad t \in (-1, +1),$$

where $\dot{u}(x, t) = tu(x, t)$. It is not difficult to prove the following lemma.

Lemma. *The equality*

$$TLu = u$$

holds.

In conclusion, from the preceding consideration, the following theorem follows.

Main Theorem. *The equation*

$$Lu = f$$

is solvable if and only if $f \in D^$ satisfies the conditions*

$$\int_{a-1}^{b+1} t\tau_{z_k}^*(x, t)f(x, t)dt dx = 0, \quad z_k \in \mathfrak{R}.$$

If these conditions are fulfilled, then the solution is unique and it can be expressed by the formula

$$u = Tf.$$

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